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Boundary layer new particle formation over East Antarctic sea ice - possible Hg-driven nucleation?


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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

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Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



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15, 19477–19536, 2015

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Abstract

Aerosol observations above the Southern Ocean and Antarctic sea ice are scarce. Measurements of aerosols and atmospheric composition were made in East Antarctic pack ice on-board the Australian icebreaker *Aurora Australis* during the spring of 2012.

5 One particle formation event was observed during the 32 days of observations. This event occurred on the only day to exhibit extended periods of global irradiance in excess of 600 W m^{-2} . Within the single air-mass influencing the measurements, number concentrations of particles larger than 3 nm (CN_3) reached almost 7700 cm^{-3} within a few hours of clouds clearing, and grew at rates of 5.6 nm h^{-1} . Formation rates of
10 3 nm particles were in the range of those measured at other Antarctic locations at $0.2\text{--}1.1 \pm 0.1 \text{ cm}^{-3} \text{ s}^{-1}$. Our investigations into the nucleation chemistry found that there were insufficient precursor concentrations for known halogen or organic chemistry to explain the nucleation event. Modelling studies utilising known sulfuric acid nucleation schemes could not simultaneously reproduce both particle formation or growth rates.
15 Surprising correlations with Total Gaseous Mercury (TGM) were found that, together with other data, suggest a mercury driven photochemical nucleation mechanism may be responsible for aerosol nucleation. Given the very low vapour pressures of the mercury species involved, this nucleation chemistry is likely only possible where pre-existing aerosol concentrations are low and both TGM concentrations and solar radiation levels are relatively high ($\sim 1.5 \text{ ng m}^{-3}$ and $\geq 600 \text{ W m}^{-2}$, respectively), such as those observed in the Antarctic sea ice boundary layer in this study or in the global
20 free-troposphere, particularly in the Northern Hemisphere.

1 Introduction

25 Nucleation in the atmosphere is important for the formation of new aerosol particles which after growth, can affect the Earth's radiative balance both directly and indirectly through their action as cloud condensation nuclei (CCN). The interaction between CCN

ACPD

15, 19477–19536, 2015

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



and radiative forcing is currently one of the largest uncertainties inherent in our understanding of the global climate (IPCC, 2013). A recent study by Carslaw et al. (2013) found that the biggest gains in reducing this uncertainty will be achieved through the study of pristine natural aerosols.

5 The Antarctic and Southern Ocean regions are among the most under-studied locations on the planet, largely due to logistical difficulties posed by their remoteness and extreme conditions. In this region, regular measurements are restricted to mid-latitude stations on continents surrounding the Southern Ocean (e.g. Cape Grim, Australia), stations on sub-Antarctic islands (e.g. Macquarie Island) or Antarctic stations. Aerosol
10 measurements in the Antarctic sea ice are particularly sparse, with only two measurements reported in the literature, both occurring in the Weddell Sea in the West Antarctic sector (Davison et al., 1996; Atkinson et al., 2012). There have been no reported aerosol measurements in the vast East Antarctic sea ice region, and as such, no characterisation of the loading there other than from model studies which, at most, have
15 been validated using data from continental stations. Recent results from this campaign (Humphries et al., 2015) have found that the Antarctic sea ice region is atmospherically distinct from both the adjacent continental and Southern Ocean atmospheres.

Aerosol nucleation is ubiquitous throughout the atmosphere, with burst events occurring in the boundary layer known as new particle formation (NPF) events. In the continental boundary layer where precursor sources are significant, NPF events are relatively common (Kulmala et al., 2004). In the marine boundary layer (MBL) however, NPF-precursor source strengths are lower and significant background aerosol populations exist (e.g. wind-produced sea-salt) which scavenge precursor vapours, preventing gas-to-particle nucleation. Consequently, only a handful of NPF events have been
20 observed in the remote MBL (e.g. Covert et al., 1996; Koponen et al., 2002; Heintzenberg et al., 2004). The Antarctic sea ice region has many characteristics of the remote MBL, however the concentration of background wind-produced aerosol populations is significantly reduced due to the ice barrier between the ocean and atmosphere. Consequently, NPF events in the Antarctic sea ice boundary layer could be more likely than

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



location, speed and direction (GPS), pitch and roll of the ship and sea temperature. Full data and meta-data are available through Reeve (2013).

Ocean capable inlets allowed measurements of atmospheric trace gases during the ocean crossing, while logistical limitations meant aerosol measurements were only conducted in the sea ice zone for 32 days between 23 September and 25 October, 2012.

2.2 Aerosol measurements

In-situ number concentrations were measured using two condensation particle counters (CPCs) every second during the voyage time within the Antarctic sea ice zone (Humphries et al., 2014). Measurements of two different size ranges were made: particles with diameters larger than 3 nm (CN₃; Model 3025A, TSI, Shoreview, MN, USA); and those larger than 10 nm (CN₁₀; Model 3772, TSI, Shoreview, MN, USA). Operational and logistical constraints meant a comprehensive aerosol measurement suite was unable to be deployed.

Air was sampled from a 3 m high mast located on the starboard side of the fore-castle (totalling approximately 12 m a.s.l.) at a rate of 130 L min⁻¹ (to minimize loss via diffusion) through 19 m of 55 mm diameter antistatic tubing. An additional 2.7 m of $\frac{1}{4}$ " conductive tubing connected into the laboratory and split the flow to both CPCs. The long inlet (used for logistical reasons) led to losses that were calibrated post-voyage against a laboratory reference inlet (designed for minimal loss) to yield a relative transmission efficiency of 0.89. The absolute losses were also calculated using two aerosol loss calculators: the first described by von der Weiden et al. (2009), while the second was written by Baron, P.A. and is available online (<http://aerosols.wustl.edu/AAARworkshop08/software/AEROCALC-11-3-03.xls>). Both calculators give similar results. For particles above 10 nm, calculated absolute inlet efficiencies ranged from 0.8 to 0.99. Consequently, the experimentally derived transmission efficiency of 0.89 was deemed suitable for this size range and applied to the data. For sizes below 10 nm, calculated inlet efficiencies decreased rapidly from around 0.8 at 10 nm, down to 0.3 at

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

3 nm. However these calculations are not valid in the laminar-turbulent flow transition regime in which this inlet system lies. Additionally, results of inlet characterisation experiments showed significant variations that led to unreliable calibration factors in this size range. Consequently, application of these calculated inlet efficiencies would be unlikely to result in more accurate data due to the high uncertainty in this size range. Instead, the same factor of 0.89 was applied to this size range and hence the results presented here are considered a lower bound for actual number concentrations.

Aerosol data were filtered to remove the significant influence of ship exhaust to achieve a dataset reflective of background aerosol loading. Filtering was performed by removal of data when relative wind directions were between 150 and 270° (assuming 0° is the bow of the ship). The influence of recirculated ship exhaust was also considered by analysis of ozone (O₃) concentrations (Klekociuk et al., 2015) and back-trajectories. The absence of any significant decreases in O₃, which is titrated out by NO_x in the ship exhaust, confirmed that recirculation did not affect our data set.

CN₃ and CN₁₀ data can be used to calculate other parameters that are useful for data interpretation. Taking the difference between these two measurements resulted in the number concentration of nanoparticles (3–10 nm; CN_{3–10}) recently formed from gas-to-particle conversion. The parameter (CN₃ – CN₁₀)/CN₁₀ has been shown before (Warren and Seinfeld, 1985; Covert et al., 1992) to allow new particle formation to be easily distinguished from background concentrations and has also been calculated for the current dataset.

Although the size information available from these measurements is basic, it is sufficient to estimate the growth rate using established equations (Kulmala et al., 2004). These equations result in good approximations under the assumption that the calculation period uses data measured within a single air-mass – a good assumption in this case, as will be demonstrated later. Using this calculation, growth is found to occur in situ if changes in the number concentration of particles larger than 10 nm were delayed relative to changes in the number concentration of 3–10 nm particles.

2.3 Trace gas measurements

Measurements of various trace gases occurred during both ocean and sea ice sections of the voyage. Sample inlets for these measurements were located fore of the exhaust at approximately 18 m a.s.l. This position meant that winds between 60 and 190° (assuming 0° is the bow of the ship) sampled ship exhaust directly. These values were filtered out using wind direction and wind speed (< 5 knots) data.

One minute averages of in situ O₃ (Schofield et al., 2014b) were measured with a dual cell ultraviolet ozone analyser (Thermoelectron 49C), sampling through a particle filtered 30 m length of $\frac{1}{4}$ " Teflon tube. The instrument was calibrated to a traceable ozone standard before and after the voyage, and zeroed weekly during the voyage using an inline O₃ scrubber for 30 min.

Six halocarbons including CH₃I, C₂Cl₄, CH₃CCl₃, CHCl₃, CH₂Br₂ (coeluted with 10–30% CHBrCl₂, Robinson et al., 2014) and CHBr₃, were measured (Schofield et al., 2014e) using a gas chromatograph with electron capture detector (GC-ECD, μ Dirac, Gostlow et al., 2010) via a particle filtered 60 m length of $\frac{1}{4}$ " Teflon tube at a rate of 1 L min⁻¹. Helium (99.995 % purity) and N₂ (99.998 % purity) were used as carrier and detector make-up gases. The GC-ECD analysed approximately 20 samples per day. Blanks and calibration chromatograms were run after every 4 sample chromatograms. The calibration air was drawn from a cylinder of clean natural air, previously filled at Niwot Ridge (Colorado, USA) at an elevation of 3.5 km a.s.l. and enriched as needed with known concentrations of the target compounds. The cylinder was supplied by the Earth System Research Laboratory (ESRL) within the National Oceanic and Atmospheric Administration (NOAA) so that the reported mixing ratios are linked directly to the NOAA halocarbon calibration scales.

Total Gaseous Mercury (TGM) was measured (Schofield et al., 2014c) at 5 min intervals using a gold trap mercury analyser (Tekran Model 2537, Tekran Instruments Corporation, Toronto, Canada¹). The inlet consisted of an inline particulate filter housed

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Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



ated at heights of 10, 500 and 1000 ma.s.l. at the ship's location during the time surrounding the 18 October particle formation event. Each calculation estimated hourly three-dimensional air-parcel locations within the specified time-frame.

At the high latitudes of the Antarctic, reanalysis datasets used to calculate trajectories rely on sparse meteorological measurements, resulting in high uncertainties, particularly when run for multiple days. Trajectories were therefore restricted to 72 h, which limited the resulting uncertainty.

2.5 Hydroxyl radical, OH

OH[•] concentration on the 18 October was calculated from values of $J(\text{O}^1\text{D})$ using the method described by Creasey et al. (2003). $J(\text{O}^1\text{D})$ was not measured directly, however using the relationship modified from Wilson (2014), it can be calculated such that:

$$J(\text{O}^1\text{D}) = 2 \left(\sum_i A_i \exp \left(-\frac{B_i}{\cos \theta} \right) \right) \cdot \left(\frac{\text{O}_3^{\text{sat}}}{300} \right)^{-\text{RAF}} \quad (1)$$

where θ is the solar zenith angle, O_3^{sat} is the total ozone column retrieved from satellite for the measurement day, and A_i , B_i and RAF (the Radiation Amplification Factor) are derived experimentally and are given in the literature cited above. The equation given by Wilson (2014) is suitable for low albedo regions, such as the Southern Ocean, however over the dense pack ice, the high albedo surface increases the actinic flux significantly. The multiplication factor of two leading the equation is used to roughly correct for this elevated albedo. O_3^{sat} was derived from OMI satellite measurements (<http://www.esrl.noaa.gov/>) to be 441 Dobson units during the event day. Under a clear sky assumption, this resulted in a $J(\text{O}^1\text{D})$ of $1.2 \times 10^{-5} \text{ s}^{-1}$. Consideration of the reaction chemistry (e.g. primary sink mechanisms using $[\text{CO}] = 60 \text{ ppbv}$ and $[\text{CH}_4] = 1.8 \text{ ppmv}$ [<http://ds.data.jma.go.jp/gmd/wdcgg/>]), the midday OH[•] concentration can be estimated to be $3.4 \times 10^6 \text{ molec cm}^{-3}$. This is in good agreement with estimates calculated using the relationship described by Rohrer and Berresheim (2006).

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



This high value is driven primarily by the high surface albedo. Direct measurements (as opposed to the estimates made here) around the Antarctic region (both continental, Mauldin et al., 2001 and coastal, Bloss et al., 2007; Kukui et al., 2014) report similarly high OH concentrations, giving some confidence that our calculations are within realistic ranges.

2.6 Box modelling

The nucleation and growth chemistry of the NPF event was investigated using a box-model version of the Two-Moment Aerosol Sectional (TOMAS) microphysics algorithm (Adams and Seinfeld, 2002; Pierce and Adams, 2009a, b). Described in detail by Chang et al. (2011), this version of TOMAS is configured to simulate the number and mass of particles by nucleation, condensation and coagulation processes within 44 log-normally spaced size bins spanning dry diameters between 0.5 nm and 10 μm . TOMAS currently includes gas phase sulfur chemistry (dimethyl sulfide [DMS], SO_2 and H_2SO_4) based on work by Chin et al. (1996), which has been shown to work well compared to other schemes in the Arctic (Karl et al., 2007). Given the similar polar environmental conditions the mechanism should be suitable for the Antarctic environment of this study.

In this environment, the model only considers DMS oxidation as a source of sulfate aerosol. This is reasonable because of the remoteness of the location from other significant sulfate sources, such as volcanoes and anthropogenic activities which are associated with high SO_2 emissions. Numerous publications (e.g. Curran and Jones, 2000; Graf et al., 2010; Shirsat and Graf, 2009, and references therein) also show that the natural DMS emissions are the major sulfur source in the region, with a Southern Ocean flux of 2 Tg S yr^{-1} compared to 16 Gg S yr^{-1} from anthropogenic and volcanic activity in the region. The short lifetime of SO_2 , at around a day (Lee et al., 2011), also suggests that transport from these distant sources in sufficient concentrations to contribute significantly to this nucleation event is unlikely and can therefore be ignored.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Recent work has posited the importance of methanesulfonic acid ($\text{CH}_3\text{SO}_2\text{OH}$; MSA) in nucleation and condensational growth processes, particularly in the marine environment (Bzdek et al., 2011; Dall'Osto et al., 2012; Dawson et al., 2012; Bork et al., 2014). A simplified inclusion of MSA into TOMAS was performed such that MSA produced from DMS oxidation was identified as H_2SO_4 in the model, and added to the condensing H_2SO_4 reservoir, essentially increasing the nucleation and condensational growth rates. This is a reasonable parameterisation for an upper bound of the effects of MSA. The growth of particles larger than 3 nm by MSA should be equivalent to H_2SO_4 as both are essentially non-volatile when condensing to particles of this size (Davis et al., 1998). This simple parameterisation is sufficient for the aims of this study and a full parameterisation of MSA nucleation and growth mechanisms should be performed as part of future work. Consequently, the sensitivity of the model to included MSA parameterisations was not quantified. However, results from this inclusion, although increasing growth rates, did not change the overall conclusions associated with these modelling results.

Two different nucleation mechanisms are available in the TOMAS model that were utilised for this study. The first, known as the empirical mechanism, is a simple activation nucleation parameterization to predict nucleation rates such that $J_{\text{nuc}} = A[\text{H}_2\text{SO}_4]$, where J_{nuc} is the nucleation rate and A is an empirical parameter (Sihto et al., 2006). This linear dependence is based on the theory that stable, but sub-critical clusters exist that require the addition of a single sulfuric acid molecule to reach critical size where stable condensational growth can occur. This mechanism, although simplified, allows easy scaling of nucleation rates. The second mechanism calculates rates and critical cluster properties using the ion-mediated nucleation mechanism (Yu and Turco, 2000; Yu, 2006, 2010). This mechanism presumes that charged molecular clusters condense around natural air ions, resulting in nucleation and growth rates significantly faster than neutral counterparts, and can produce nucleation in situations where traditional nucleation mechanisms are unfavourable. Other sulfuric-acid based schemes (e.g. binary or kinetic nucleation) were not tested because they either produce similar (kinetic) or

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

smaller (binary) nucleation rates than the ones tested here (Yu et al., 2010). Although the model employed does not explicitly simulate schemes of sulfuric acid with ammonia, amines or organics, development to include these schemes was not deemed necessary for this study because of the low concentrations that are likely in the East Antarctic sea ice region where we were far from the region's biggest (although still small in the global context) sources of biological activity (see below) and the continent (e.g. Legrand et al., 1998; Davison et al., 1996; Ayers and Gras, 1980; Gras, 1983).

A range of typical input variables to the TOMAS box model, shown in Table 1, was used to account for uncertainties in the nucleation rate. The model was initialised with a pre-existing aerosol size distribution (single or dual mode distributions with variable modal median diameter and number concentration, with a fixed width, σ , of 2 which is unit-less in log-normal space) and concentrations of DMS and OH[•] that spanned the possible values in the field. Simulations were performed with A factors ranging from 10^{-10} , suitable for clean polar environments, to 10^{-6} , those found to suit continental locations (Chang et al., 2011). The pre-existing aerosol is assumed to have the density and hygroscopic growth properties of ammonium bisulfate – an assumption with minimal effect compared to changes in size distribution (Chang et al., 2011). The box presumes a pulse of DMS emission, followed by zero emissions. This is unlikely to be representative of measurements and is one of the constraints/uncertainties given careful consideration in the interpretation of results. Initial concentrations of MSA, SO₂ and H₂SO₄ are assumed to be zero, and MSA and H₂SO₄ are the only included condensable vapours (i.e. no species such as organics) and are assumed to have the same condensable properties. Three estimates of MSA fraction produced from DMS oxidation were used including 0.1, 0.3 and 0.5 (Chen et al., 2012). OH[•] was kept constant throughout the simulation, with concentrations chosen to span likely values, as well as extremes for model validation. Cloud cycling is not accounted for in the model. Simulations were run for 24 h, which spans more than the total time a constant air-mass was present over the measurement site. Over 1000 simulations were run, including all

permutations of model inputs, inclusion of MSA condensation and the two nucleation mechanisms.

3 Results

The NPF event will be investigated in this paper in the following manner. Trends in aerosol data are first described to give a general overview. After determining the air-mass influences throughout the period, the formation and growth rates, as well as the loss mechanisms and rate, are determined. The chemistry of nucleation is then investigated. This is performed first by investigating correlations of aerosol data with photochemical indicators, after-which the chemistry is assessed to determine whether known mechanisms can explain observations.

Figure 2 shows a time series of aerosol parameters CN_{3-10} , CN_{10} , and the $(CN_{3-10} - CN_{10}) / CN_{10}$ ratio for a week of the voyage encompassing the NPF event on the 18 October, 2012. Particle concentrations on this day are observed to be significantly higher than the background concentrations experienced throughout the remainder of the voyage. Background CN_3 concentrations had a maximum of 2800 cm^{-3} (median and mean of 767 and 939 cm^{-3} , respectively) and a ratio value that varies minimally around two (described further in a future publication, Humphries et al., 2015). The particle formation event on the 18 October was the only one during the full 32 day measurement period that showed both rapid increases in particle concentrations ($> 1000 \text{ cm}^{-3} \text{ h}^{-1}$) beyond background values and evidence of growth. One other high concentration period was observed but due to the absence of growth indicators, as well as back-trajectory analyses, this was attributed to variations in the age of air-masses which supply background aerosol concentrations to the region. The ratio $(CN_{3-10} - CN_{10}) / CN_{10}$ showed very little variation throughout the record (reaching a maximum, not shown, of 4.8 at any other time), but increased by almost an order of magnitude near the start of the event to a maximum of over 16.4, decreasing as particles grew to diameters larger than 10 nm. This corroborated the attribution of the event as in situ formation rather than fluctua-

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

tions in background populations. A more detailed view of the formation event is plotted in Fig. 3. At the peak of the event, CN_{3-10} number concentrations reached above 6100 cm^{-3} within a few hours, while CN_{10} concentrations reached above 1600 cm^{-3} ($\text{CN}_3 \approx 7700 \text{ cm}^{-3}$). A breakdown of the event is given as follows, with numbering referenced in Fig. 3. An air-mass change, at the beginning of period I, resulted in CN_3 number concentrations dropping from ~ 2000 to 390 cm^{-3} just prior to the event, well below average. After an overcast morning, clouds cleared at II, providing the highest levels of solar radiation observed throughout the 32 day measurement period. At III, number concentrations of CN_{3-10} began to increase rapidly, with no simultaneous increase in CN_{10} (remaining at background values of 120 cm^{-3}), suggesting nucleation. About 1.25 h later, period V began with growth into the CN_{10} size range, increasing their number concentrations and simultaneously decreasing the rate of increase of CN_{3-10} . Further increases in CN_3 number concentrations ceased at the beginning of period V, likely the result of nucleation rates slowing (due to decreased solar radiation) to the same rate as growth into the CN_{10} size bin (observed from the continued increase in CN_{10} during period V). At the beginning of VI, particle concentrations in both size bins began to decrease (likely the result of coagulation together with reduced nucleation) coincident with zero solar radiation. The rate of decline significantly increased during period VII due to an air-mass change (determined below), ending in background concentrations approximately eight hours after nucleation ceased. Variability in number concentrations, in particular the highlighted peaks, are suggested to be fluctuations in background concentration due to their simultaneous variation in both size bins and their subsequent absence in ratio data.

3.1 Identifying a single air-mass

To understand the changes in aerosol data, the various influences on the air-mass must be characterised to determine possible sources, sinks and other changes. This was

done by consideration of sea ice conditions, measured in situ meteorology, calculated back-trajectories, and ancillary composition data such as O₃ and TGM.

During the days surrounding the formation event, the ship was stuck, drifting with the pack ice due to weather patterns that caused convergence of the ice floes. This resulted in a lead-free region that was determined by helicopter survey to stretch for at least 20 km in every direction. This expanse of homogeneous, low profile (and generally multi-metre thick) sea ice suggests minimal source contributions from the (ice or ocean) surface since fluxes through the ice are low (e.g. Nomura et al., 2010) and open ocean leads were rare or non-existent.

A selection of meteorological data is included in Fig. 3. Just prior to the event, relative humidity (RH) decreased from around 96 to 85 %, while air temperature increased slightly from -8 to -7 °C, with a 2 °C variance around the higher temperature. This pattern of decreased RH and slightly increased temperature is consistent with other measurements of NPF in the MBL (Covert et al., 1992). Although wind direction was consistently from the east in the days preceding, during and after the event, changes in RH and air temperature began simultaneously with the drop in aerosol number concentrations, suggesting a change in air-mass. Wind speed was below 7 ms⁻¹ throughout the event, with high wind speeds up to 20 ms⁻¹ 24 h both before the event, consistent with the passage of a low pressure system (967 hPa), and after, corresponding to the passage of a high pressure ridge (1000 hPa). Although these calm conditions could suggest ship emissions were being sampled, the lack of anthropogenic signatures in other data (e.g. O₃ or high frequency particle concentration fluctuations) indicated background air was measured. Other measurements (Raes, 1995; Clarke et al., 1998) found similar low wind speed conditions were necessary for nucleation in the MBL due to lower pre-existing aerosol surface area because of reduced wind-induced sea-spray aerosol. Sea-spray aerosol is unlikely to be as important in the sea ice region because of reduced exposed ocean, however aerosols from blowing snow (Yang et al., 2008) and frost-flowers (Rankin et al., 2002) could occur in higher winds. These

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



aerosols would scavenge vapours and prevent nucleation, making the low wind speed similarly important for low pre-existing aerosol numbers in this environment.

Figure 4 shows back-trajectories terminating at 10 m above mean sea level for the period spanning the formation event. Although in situ wind direction was consistently from the east during this period, 72 h trajectories show significant changes in air-mass history. During the high aerosol concentrations prior to the event, easterly winds brought air that had been in the free-troposphere only 36 h prior to measurement that was representative of background populations (Humphries et al., 2015). The air-mass change identified as period I of Fig. 3 and the first red track of Fig. 4a, shows the first signs of air changing direction within the 12 h prior to the event. The lower aerosol concentrations just prior to the event were found to correspond to air-masses that had recently come from the region north-west of the measurement location. During the period of increasing CN (Fig. 4b), trajectories were found to “recirculate” around this north-west region for about 24 h, before being measured. This gives a meso-scale (~ 500 km) spatial bound to the chemistry and conditions leading to the formation event, rather than the local-scale. A change from this recirculation occurs at approximately 15:00 (UTC; Fig. 4c), when the decreasing aerosol concentrations began to plateau before a sudden decrease, caused by an abrupt change in air-mass history that increased influence from areas further west of the recirculation region. Overall, trajectories support the conclusion that a single air-mass influenced the measurement location throughout the period of interest, and that this air-mass came from the north-west.

Independent measurements of atmospheric components, such as O_3 , can help in identifying air-mass influence. In Fig. 5a, O_3 can be seen to be relatively constant throughout the nucleation period at approximately 27.5 ppb, with abrupt changes occurring at times previously identified (via CN and trajectory data) as air-mass changes. The constant O_3 concentrations between these abrupt changes then, are reflective of the influence of a single air-mass.

TGM is another independently measured atmospheric component that shows (Fig. 5a) significant and simultaneous change with the air-mass changes identified

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

previously. Notwithstanding the recent establishment of a number of continuous monitoring stations, background TGM concentrations are still relatively rare in the Southern Hemisphere (Slemr et al., 2015). Despite short term spikes above 1.5 ng m^{-3} , our measurements generally showed background values ranging from 0.7 to 1.1 ng m^{-3} , similar to the most recent Southern Hemisphere values that range between about 0.8 and 1.1 ng m^{-3} (Slemr et al., 2011, 2015). During the period being discussed, concentrations increased from less than 1.0 ng m^{-3} , up to 1.5 ng m^{-3} , remaining at this concentration for a sustained period relative to previous spikes, before returning back to background values less than 1 ng m^{-3} , coincident with the ending air-mass change. This corroborates the previous inference of a single air-mass identified to be influencing the event.

If nucleation is occurring at the meso-scale, steady or slowly changing aerosol trends will be observed in particle concentration. Conversely, if events are occurring over the scale of 10s of metres, sudden and large fluctuations in number concentration would be expected on short timescales on the order of seconds to minutes. The generally steady trends in aerosol number concentrations during this event support the conclusion of a meso-scale homogeneous air-mass.

Since a single air-mass is influencing the measurement site, it can be concluded that the physics and chemistry over this meso-scale are homogeneous. Consequently, during this NPF event, conclusions about the growth of aerosols and changes in atmospheric composition from measurements from our stationary site can be made confidently (assuming horizontal homogeneity) without worrying that variability is a result of advective effects due to changing air-mass.

3.2 Formation and growth rates

Formation rates of 3 nm particles, J_3 , averaged over a given time interval, were calculated using the method described by Kulmala et al. (2004). Assuming the effects of coagulation and transport were small during the chosen time period, a reasonable assumption when air-masses are relatively clean and homogeneous, as found in this

in-situ processes, such as coagulation, are more likely to be driving the declining number concentrations than other processes determined by air-mass changes, such as horizontal dispersion and advection.

The influence of vertical mixing and coagulation can be modelled assuming a well-mixed boundary layer capped with an inversion where particle free air is being entrained from the free troposphere. This can be quantified by the following equation (Davison et al., 1996) which describes the time rate of change of CN concentration:

$$\frac{dn}{dt} = -n \left(\frac{kn}{2} + \frac{v_d}{H} \right) \quad (2)$$

where n is the total CN concentration, t is time, k is the coagulation coefficient, v_d is the entrainment velocity and H is the boundary layer mixing height. Integration yields a function dependent only on k , v_d , H and n_0 , the initial concentration, such that

$$n = \frac{ace^{-bt}}{1 - ce^{-bt}} \quad (3)$$

where $b = v_d/H$, $a = 2b/k$ and $c = n_0/(a + n_0)$. Taking an entrainment velocity (v_d) of 0.34 cm s^{-1} (from calculations performed with Weddell Sea measurements by Davison et al., 1996), a calculated boundary layer height (H) of 200 m (based on trajectories calculated with the UK Met Office's Numerical Atmospheric-dispersion Modelling Environment, Schofield et al., 2014a) and an initial CN concentration of 7500 cm^{-3} , a coagulation coefficient (k) of $(8 \pm 2) \times 10^{-9} \text{ cm}^{-3} \text{ s}^{-1}$ results in the best fit to period VI of the data, shown in Fig. 6. The two variables v_d and H , are assumed from the literature and need an assessment of sensitivity. Calculation of k is insensitive to boundary layer height, but highly sensitive to entrainment velocity. Reducing the entrainment velocity by a factor of 10 increases k by only a factor of two and results in a similarly good fit. However, an increase in the entrainment velocity by an order of magnitude results in the inability to find a value of k that fits the measurements, a fact consistent with the unrealistic nature of this change. Although the uncertainty around the entrainment

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



velocity parameter is high, it is still possible to conclude that coagulation and vertical mixing are sufficient to explain the loss.

The reduction of CN concentrations during the event is therefore likely to be explained by three periods. The first, period V, involves very slow loss during which photochemical production of CN is still possible before sunset. During this period, coagulation has already started but CN concentrations are being maintained by continued production. At sunset, production ceases and coagulation dominates the changing concentrations of period VI. Later, CN losses increase drastically during period VII, coinciding with the air-mass change described previously and shown in Fig. 4c.

Comparison of this coagulation coefficient ($(8 \pm 2) \times 10^{-9} \text{ cm}^{-3} \text{ s}^{-1}$), derived during period VI, to those of coefficients calculated using Fuchs formula at -11°C (Seinfeld and Pandis, 2006), yields information about the likely sizes of coagulating aerosols. Fuchs formula calculates the coagulation coefficient by considering the coagulation of two particles with wet diameters D_1 and D_2 . Resulting calculations yield coagulation coefficients of $(8 \pm 0.1) \times 10^{-9} \text{ cm}^{-3} \text{ s}^{-1}$, only when a D_1/D_2 ratio of 5.1 ± 0.3 exists. Ratio values above 5.5 or below 4.6 correspond to coagulation coefficients well outside the range plausible for a reasonable fit to the data. Particles with the appropriate diameter ratio, for example 3 and 17 nm, or 4 and 21 nm, are therefore the most likely to be present and responsible for the observed coagulation rate.

3.4 Aerosol surface area

If existing aerosol concentrations are high enough to provide a sufficient surface area, precursors are more likely to condense onto these surfaces rather than nucleate to form new nuclei. Although aerosol surface area cannot be calculated accurately from the available data, approximate values can be inferred using Antarctic aerosol size distributions from the literature, along with size information revealed from theoretical calculations (above). Assuming a mean background aerosol population diameter of 100 nm, consistent with measurements by Koponen et al. (2002), the air-mass change prior to the event resulted in a transition in pre-existing aerosol surface area from 19 to

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



(a proxy for biological activity), could be compared to previous studies relating chlorophyll to SOA formation. Studies by Hu et al. (2013) showed that chlorophyll concentrations necessary for organic concentrations sufficient for SOA formations are of the order of $10\text{--}60\text{ mg m}^{-3}$, which were a result of summer phytoplankton blooms in Prydz Bay, Antarctica. During SIPEXII, no such phytoplankton blooms were observed, and chlorophyll concentrations reached a maximum of 0.4 mg m^{-3} (median of 0.15 mg m^{-3}) during the aerosol measurement period, with order of magnitude lower concentrations on the event day, averaging $0.06 \pm 0.03(2\sigma)\text{ mg m}^{-3}$. Chlorophyll *a* concentrations acquired from the NASA Oceancolor website (<http://oceancolor.gsfc.nasa.gov>) confirm in situ data, with concentrations well below 0.3 mg m^{-3} in the hundreds of kilometers surrounding the measurement location. Despite this evidence, nucleating species could be produced in situ. However, the measurements taken aboard SIPEXII strongly indicate that their concentrations, if present, are at a minimum during the event compared to other periods of the voyage, suggesting their involvement is unlikely.

3.6.2 Halogen chemistry

Active nucleation chemistry involving iodine oxidation has been demonstrated to occur rapidly on the local scale when coastal kelp-beds are exposed during low-tide periods (O'Dowd et al., 1998, 1999, 2002a). This chemistry is found to coincide with high levels of iodine monoxide (IO) of up to 6 pptv (Saiz-Lopez and Plane, 2004). In-situ (Frieß et al., 2001, 2010; Saiz-Lopez et al., 2007b; Atkinson et al., 2012) and satellite (Saiz-Lopez et al., 2007a; Schönhardt et al., 2008, 2012) measurements of IO (and I_2) reveal Antarctica as a global iodine hotspot, particularly in the West Antarctic during spring, where concentrations of 6 pptv are frequent, and midday values up to 20 pptv have been observed. Although these high concentrations suggest that nucleation from iodine chemistry could be prolific in the Antarctic region in general, measurements during this voyage found IO concentrations average below about 0.5 ppt, and during the nucleation event, concentrations were at the detection limit of approximately 0.1 ppt. The influence of a single regional air-mass, combined with the exceedingly low IO

concentrations measured during the event, make it unlikely that iodine chemistry was responsible for this observed nucleation event.

3.6.3 Sulfur chemistry

Nucleation mechanisms involving sulfuric acid (e.g. binary nucleation of H_2SO_4 and H_2O) are the most common in the atmosphere. Because of this ubiquity, these mechanisms are often the only ones included in chemistry and climate models. In this study, it was not possible to determine whether sulfur chemistry was involved in nucleation using the available measurements. To fill this gap in measurements, the TOMAS box model was employed to examine whether existing nucleation schemes involving sulfuric acid could explain the observed nucleation rates. The source of H_2SO_4 employed in this box model was the biggest natural sulfur source, DMS. Utilisation of a box model was suitable for this study primarily because of the single homogeneous air-mass influencing the measurement location during the event.

The purpose of the modelling study was to determine whether existing H_2SO_4 -dependent nucleation schemes could account for observations under the measured conditions. If, even under the most favourable conditions for DMS driven nucleation and growth (high DMS and OH^\cdot concentrations), the model cannot produce the observed number and rate of growth, it is likely that another species or alternate chemical mechanism is driving the nucleation chemistry. Because the study aimed for a qualitative result, it was suitable to account for the (potentially high) uncertainty associated with the input variables by using a wide range of both plausible and extreme values, outlined earlier in Table 1.

Analysis of results determined that no combination of input parameters utilising either nucleation mechanism included in the model could simultaneously produce both the formation and growth rates necessary to explain the observed nucleation event. Inclusion of MSA chemistry into the model increased growth rates significantly, however this modification still did not result in any simulations that were able to reproduce observations.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

which experienced significantly high peaks of cloud-edge effect radiation. The period of declining TGM ended simultaneously with the appearance of a temporary cloud period around 06:00 (UTC), and did not begin to decline again when this cloud cleared and solar radiation had dropped below 600 W m^{-2} . This correlation suggests possible photochemically driven chemistry involving TGM that may require high incoming solar radiation levels – chemistry that is consistent with previous studies that have observed photochemically derived mercury oxidation products in the Arctic atmosphere (Lindberg et al., 2002).

This four hour period of declining TGM concentrations, from an initial value of 1.50 ± 0.05 down to $1.16 \pm 0.05 \text{ ng m}^{-3}$, represents an atmospheric TGM depletion of about 20%. The decline, along with the four hour recovery back to baseline, was most likely caused by in situ chemistry, rather than an air-mass change which has been discounted for this period by previous arguments. As discussed by Lindberg et al. (2002), photochemical reactions producing halogen and OH radicals are able to react with Hg^0 to produce oxidised compounds such as HgO , $\text{Hg}(\text{OH})_2$, HgX_2 or HgXY (where X and Y are halogens). Additionally, $\text{Hg}(\text{II})$ products are able to be photolytically released directly to the atmosphere from surfaces such as the snowpack (United Nations Environmental Programme, 2008). Oxidation products have properties that make them suitable for nucleation processes, such as being solid at equilibrium under atmospheric conditions making them unlikely to evaporate from nucleating clusters (Schroeder and Munthe, 1998). Production of these species can occur by direct emission or in situ chemistry. Ordinarily, Hg oxidation products, once produced, are readily scavenged by pre-existing aerosol populations. This interpretation could be largely driven by the location of most Hg studies being in regions of high pre-existing aerosol surface area. In more pristine conditions where pre-existing aerosol surface area is low, such as those observed in this Antarctic study, these oxidation products have the potential to nucleate rather than adhere to existing surfaces.

To further investigate the possibility of nucleation involving Hg, a lagged correlation was calculated between TGM and CN_{3-10} . Unfortunately, there was insufficient

1987) and a 1 : 1 stoichiometry exists between Hg^0 and reactive gaseous mercury products (e.g. HgO , HgBr_2 , and HgCl_2), the full four hour depletion period produces a total of $(8.4 \pm 0.2) \times 10^5 \text{ molec cm}^{-3}$. Assuming all of this contributed to the 5500 cm^{-3} of aerosol nuclei, each nucleus would have around 150 molecules. Assuming sphericity and that bulk properties are similar to aerosol properties, aerosol diameters using purely TGM molecules range from 2.2 to 2.8 nm (depending on the chemical identity). This is sufficient for aerosol genesis, provided other condensable species are available for further growth.

Oxidised mercury products are formed regularly in the atmosphere, and are found frequently in atmospheric particulate matter (termed particulate mercury, Hg(p)). The processes leading to the formation of Hg(p) generally involve the oxidation of Hg^0 to reactive gaseous mercury compounds (Hg(II)) which are of sufficiently low vapour pressure that they readily bind to pre-existing aerosol surfaces. This process of oxidation to deposition occurs on the timescale of hours to days (United Nations Environmental Programme, 2008). In the case where pre-existing aerosol concentrations are low, such as those observed in the event being described, we speculate that reactive gaseous mercury may nucleate by itself, or together with other species, to form new aerosol particles.

The day preceding the new particle formation event experienced a significant bromine-driven mercury depletion event (shown in the Appendices), similar to those commonly observed in the Arctic spring-time (e.g. Schroeder et al., 1998; Lu et al., 2001; Ebinghaus et al., 2002; Boudries and Bottenheim, 2000; Lindberg et al., 2002) together with high snowfall. The chemistry involved in this mercury depletion results in oxidised mercury species being deposited onto the snowpack surface of the sea ice from which it can enter the hydrosphere and biosphere, or can be re-emitted back to the atmosphere photolytically. Given the high solar radiation that influenced the event, it is possible that reactive mercury compounds were photolytically re-emitted directly from the surface from the previous day's deposition event. These direct emissions of

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



condensable material (which may or may not undergo further in situ chemical reactions), could be sufficient to explain observations.

It is also possible that Hg^0 , readily available from long-range transport, is oxidised in situ to Hg(II) compounds which then condense to form the new nuclei. Hg^0 can be oxidised by numerous species including O_3 (Snider et al., 2008), OH^\cdot and halogens (X), producing a range of Hg(II) compounds suitable for condensation (some listed above).

The identity of the oxidant in this event is worthwhile considering. Steady concentrations of O_3 at around 27.5 ppbv (at 1 atm and 0°C , $27.5 \text{ ppbv} \approx 7.5 \times 10^{11} \text{ molec cm}^{-3}$) throughout the period, together with a slow reaction rate with Hg^0 (gas phase rate constant of $(3 \pm 2) \times 10^{-20} \text{ cm}^3 \text{ molec}^{-1} \text{ s}^{-1}$, Hall, 1995) suggest it was not involved. OH^\cdot concentrations on the 18 October were calculated to be high, likely due to a combination of enhanced solar radiation levels and high surface albedo, at $3.4 \times 10^6 \text{ molec cm}^{-3}$, although the slow rate constant of $(9.0 \pm 1.3) \times 10^{-14} \text{ cm}^3 \text{ molec}^{-1} \text{ s}^{-1}$ (room temperature) (Pal and Ariya, 2004), suggest it is insufficient to fully explain nucleation. Discussed in detail in the Appendices, an assessment of measurements and the literature suggests that although the reaction rates of Hg with halogen radicals are high (rate constants are $(1.0 \pm 0.2) \times 10^{-11} \text{ cm}^3 \text{ molec}^{-1} \text{ s}^{-1}$ and $(3.2 \pm 0.3) \times 10^{-12} \text{ cm}^3 \text{ molec}^{-1} \text{ s}^{-1}$ for Cl^\cdot and Br radicals, respectively, Ariya et al., 2002; Stephens et al., 2012), the atmospheric halogen species, Cl^\cdot , Br^\cdot and I^\cdot , or their oxidation products, are present in insufficient concentrations to be involved as oxidants in this event.

Whatever the mechanism of production or identity of the oxidant, oxidised mercury species have the potential to condense, either homogeneously or on pre-existing sub-critical nuclei, to form critical nuclei from which rapid growth can occur via condensation of other species (or other oxidised Hg compounds). It is possible then, that Hg was necessary for aerosol nucleation in this event where background aerosol surface area was sufficiently low. A possible nucleation pathway, driven by Hg, is outlined in Fig. 9. Importantly, the oxidized Hg resulting from this reaction mechanism may be recycled back to TGM (given the trends and relationship presented in Fig. 8), resulting in no net effect on the mercury budget. Although this mechanism is theoretically possible,

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Hg driven aerosol nucleation has not been previously described, and correlations in experimental data could be coincidental, rather than reflective of a causal relationship.

4 Summary and discussion

Aerosol concentrations at two size thresholds, together with ancillary atmospheric composition data, were used to investigate the only in situ particle formation event that occurred during the 32 days of measurements during a marine science voyage in the pack ice off the East Antarctic coast. The significant particle formation event occurred within a single, homogeneous air-mass and resulted in CN₃ particle concentrations reaching almost 7700 cm⁻³ within a few hours after being at background values under 400 cm⁻³. Formation rates of 3 nm particles ranged from 0.2–1.1 ± 0.1 cm⁻³s⁻¹, higher than most events observed in the Antarctic region, and similar to a recently described rare continentally derived event (Kyrö et al., 2013). Estimated growth rates were higher than those previously observed in the region, being 5.6 ± 0.9 nm h⁻¹ during growth from 3 to 10 nm diameters. An assessment of the reduction in particle numbers found that loss processes were dominated by coagulation, with calculations yielding a coagulation coefficient of (8 ± 2) × 10⁻⁹ cm⁻³s⁻¹, consistent with coagulation of particles with diameters, D₁ and D₂, and a D₁/D₂ ratio of 5.1 ± 0.3 (e.g. 3 and 17 nm particles), corresponding to the likely sizes present.

The event was characterised by pre-existing aerosol surface areas below 5 μm² cm⁻³, along with abnormally clear cloud conditions that resulted in incoming solar radiation fluxes higher than observed on any other day during the voyage, reaching over 600 W m⁻² for four hours, and 700 W m⁻² at midday.

An evaluation of the mechanism responsible for nucleation found that known sulfur, halogen or organic chemistry could not explain observed nucleation and growth rates. Statistically significant correlations were found to exist between CN₃ and TGM concentrations, as well as incoming solar radiation. A mercury driven nucleation mechanism is proposed as a plausible explanation for the observed nucleation.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the globe (United Nations Environmental Programme, 2008), could lead to nucleation when pre-existing aerosol surface areas are low enough.

The prerequisites for nucleation via this mechanism are unlikely to be confined to the polar boundary layer. Conditions of low pre-existing aerosol surface area are common above clouds in the free-troposphere and stratosphere. Coincidentally, these locations also experience high levels of solar radiation. In the Northern Hemisphere in particular, TGM concentrations are also likely to be sufficient at these altitude levels given the significantly higher boundary layer concentrations (Schroeder and Munthe, 1998; Slemr et al., 2003). This suggests that this nucleation mechanism could be occurring throughout the atmosphere.

Other high mercury, low particle environments may also host this type of nucleation. Coal-fired power plants are the largest single source of mercury in most countries, with an estimated 25 % of global anthropogenic emissions to the atmosphere (United Nations Environmental Programme, 2008). Power plant exhaust plumes can have TGM concentrations enhanced up to 6.4 ng m^{-3} above ambient concentrations (Deeds et al., 2013) that are commonly much higher than those observed here (e.g. ambient concentrations at Alert Station in the Canadian Arctic are 2 and over 5 ng m^{-3} in Asia, United Nations Environmental Programme, 2008). Filtering of primary particulate emissions from power plant exhaust is mandatory in many countries around the world. Within these plumes, it is possible that the ratio between Hg^0 concentrations and pre-existing aerosol surface area is high enough that secondary particle formation driven by mercury is possible. If this were the case, this chemistry could lead to an additional aerosol burden in industrial environments which should be quantified.

The correlations between nucleation mode aerosol and mercury observations, and thus the proposal of this mechanism, have likely not been suggested before because coincident measurements of particle formation events and mercury are rare. These two measurements only occurred in this study serendipitously. Additionally, high aerosol concentrations in commonly sampled areas results in reactive, low-volatility mercury

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



products being scavenged by pre-existing particles, removing them from the atmosphere and preventing their nucleation.

It is important to note that the assumption used in this work, that the transmission efficiency of the aerosol sample inlet was the same across all sizes (discussed in Sect. 2.2) is flawed, particularly at sizes below 10 nm. If better characterisation of the inlet were possible at the sub-10 nm size range, it is likely that CN_{3-10} would increase significantly ($\sim 50\%$), which would increase both formation and loss rates. If this correction were applied and CN_{3-10} number concentrations increased, the problems associated with the model's inability to reproduce observations would be exacerbated and strengthen the major conclusions of this work.

Results described here represent only a basic measurement suite. More detailed measurements of aerosols, including full size distributions and chemical and elemental composition (e.g. fractions of organic matter, sea salt, sulfate, or other species such as mercury), along with detailed mercury and sulfur speciation, are necessary for full experimental characterisation of this chemistry. Further aerosol focussed studies are necessary in this environment to more completely characterise the aerosol loading as well as the frequency and impact of new particle formation. Ideally, these measurements would be both campaign based and long-term (e.g. the O-buoy system, Knepp et al., 2010), although the latter presents significant logistical challenges in the dynamic sea ice zone. Laboratory (in particular chamber studies) and modelling based studies would help to determine, both qualitatively and quantitatively, whether this chemistry has any atmospheric significance. If substantiated through further observations, the newly described chemistry could be important in regions where pre-existing aerosol populations, and resulting scavenging rates of oxidised mercury products, are low.

Appendix A: Mercury depletion on 17 October

As discussed in the main text, the high Total Gaseous Mercury (TGM) concentrations on 18 October followed a significant depletion that occurred on 17 October, as shown

19511

ACPD

15, 19477–19536, 2015

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



running the peak analysis software for the μ Dirac. N. Harris and A. Robinson would like to thank NERC for financial support through an Advanced Research Fellowship and the CAST project.

Data access. Data from the voyage are available at the Australian Antarctic Data Centre (AADC). Dataset DOIs will be made public upon publication.

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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

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R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

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R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

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R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Table 1. Range of model input parameters to test sulfur chemistry.

Parameter	Tested Values
Background aerosol concentration (cm^{-3})	100, 200, 400, 500, 600, 1000
Background Mode Size (nm)	0.5, 1, 1.5, 2, 3, 10, 30, 50, 100, 150, 300
[OH] (molec cm^{-3})	2×10^5 , 4×10^5 , 5×10^5 , 6×10^5 , 8×10^5 , 1×10^6 , 5×10^6 , 1×10^7
Initial [DMS] (pptv)	100, 200, 300, 400, 500, 600, 700, 800
A (s^{-1})	10^{-10} , 10^{-9} , 10^{-8} , 10^{-7} , 10^{-6}

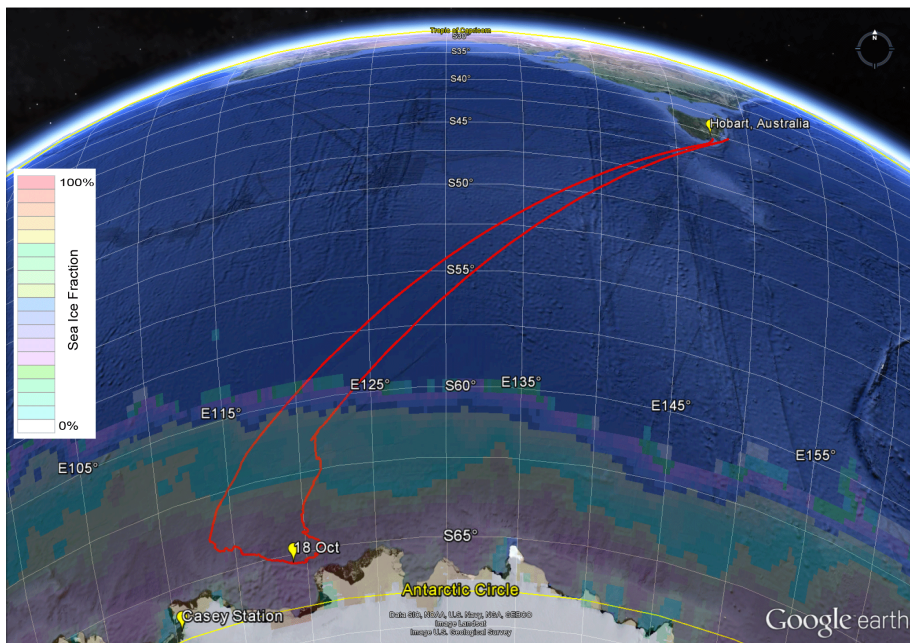


Figure 1. Ship track (red) relative to the departing port (Hobart, Australia) and the Antarctic continent. Sea ice concentration is shown in colours around the continent and is derived from Extended Reconstructed Sea Surface Temperature (ERSST) Version 3b, October 2012 mean. Imagery: Google Earth.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

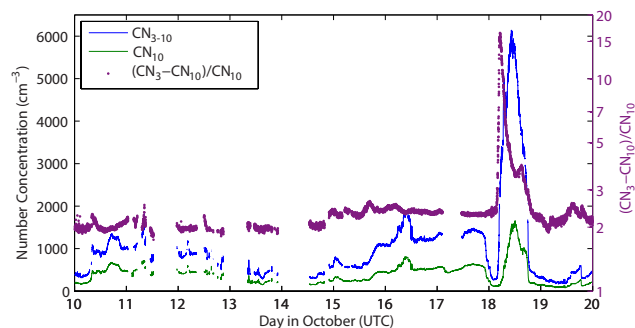


Figure 2. Time series of aerosol parameters for a 10 day subset of the 2012 SIPEXII voyage. Aerosol parameters include calibrated number concentrations of CN_{3-10} (blue), CN_{10} (green) and the ratio CN_{3-10}/CN_{10} (red). The ratio has been shown previously (Warren and Seinfeld, 1985) to be a good indicator of new particle production when limited measurements are available.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

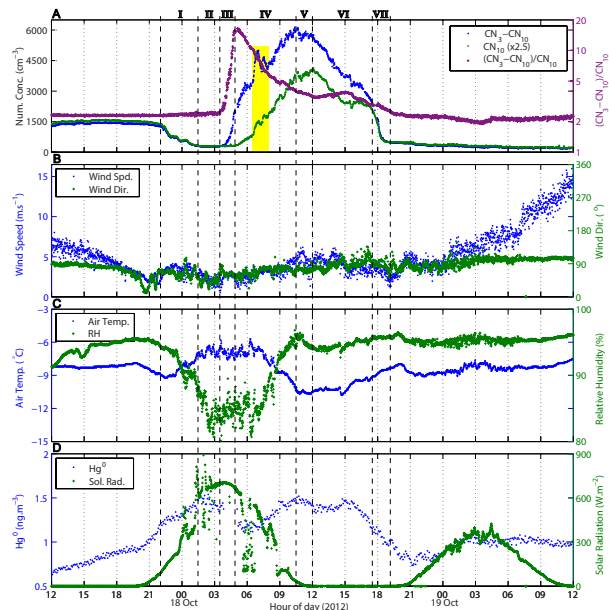


Figure 3. Detailed view of particle formation event along with selected meteorological parameters and TGM concentrations. See text for details of numbered events. **(a)** Number concentrations of CN_{3-10} , CN_{10} , and $(\text{CN}_3 - \text{CN}_{10}) / \text{CN}_{10}$ show the structure of the formation event clearly. Note the delay between increases in the two size bins, and the linearity of the increase in ratio during period III. The highlighted period shows short-term spikes in both aerosol size bins, with no signal in the ratio data, suggesting this is a result of variations in background concentrations. **(b)** In-situ wind direction during the period is constantly from the east, with minimal variation. Wind speed was low throughout the formation event, remaining below 7 ms^{-1} . This is in contrast to the days before and after the event where high wind speeds were observed. **(c)** The beginning of the event saw slight increases in temperature simultaneous with a 10% decrease in relative humidity. **(d)** Solar radiation was found to reach higher levels than any other during the voyage, and decreases in TGM were found to correspond with the period of highest solar radiation.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

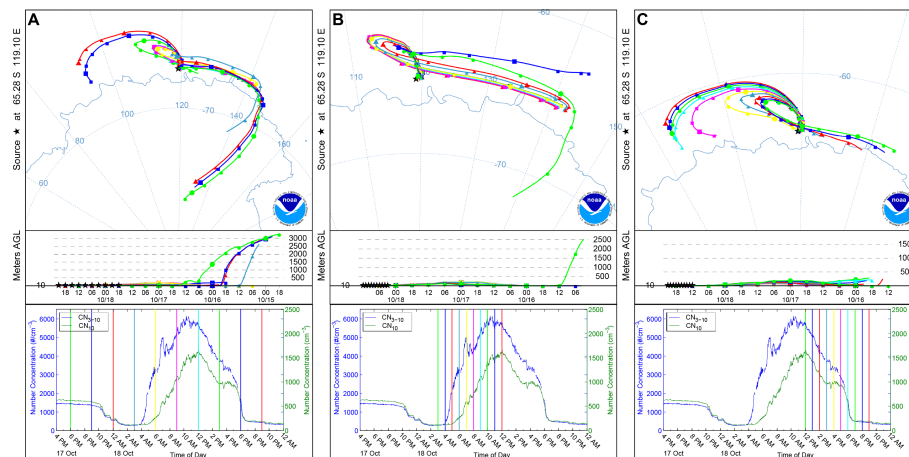


Figure 4. Back-trajectories released from 10 m (a.s.l.) calculated using NOAA's HYSPLIT model and driven by GDAS meteorological data, run for 72 h during the 18 October particle formation event. Coloured lines in the lower plots show the release time of the trajectory relative to the particle concentration. **(a)** Ten trajectories, one released every 3 h between 17 October 18:00 and 18 October 21:00, show the air-mass change from a south-east background, to a more local influence, then back to background from the west. **(b)** Released hourly and spanning the increase in particle number concentration, between 03:00 and 12:00 18 October, trajectories show that air-masses recirculate to the north-west of the ship after travelling along the sea ice surface. **(c)** Hourly trajectories during the decrease in particle number concentration, between 12:00 and 21:00 18 October, show the local impact steadily being diluted as influence from further west increases. The fastest drop in aerosol concentration is found to occur when recirculation ceases and westerly influence dominates.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

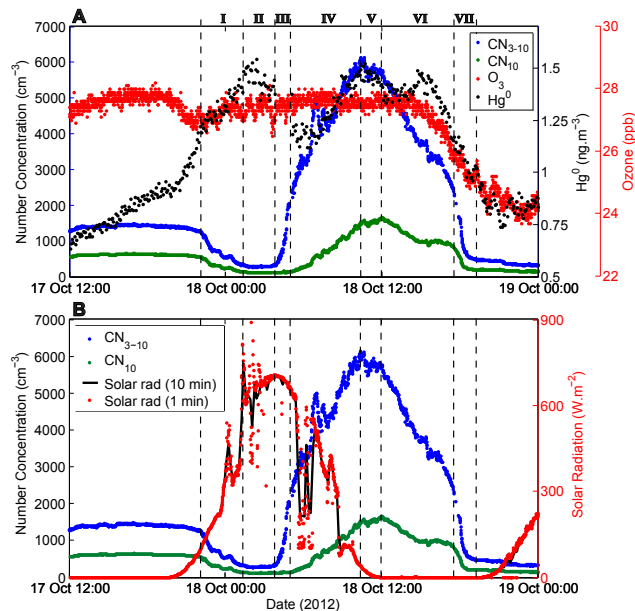


Figure 5. Close up time series showing how various other parameters vary relative to the particle formation event. Aerosol concentrations are included for reference. **(a)** O₃ and TGM. **(b)** Solar radiation (1 and 10 min averages).

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

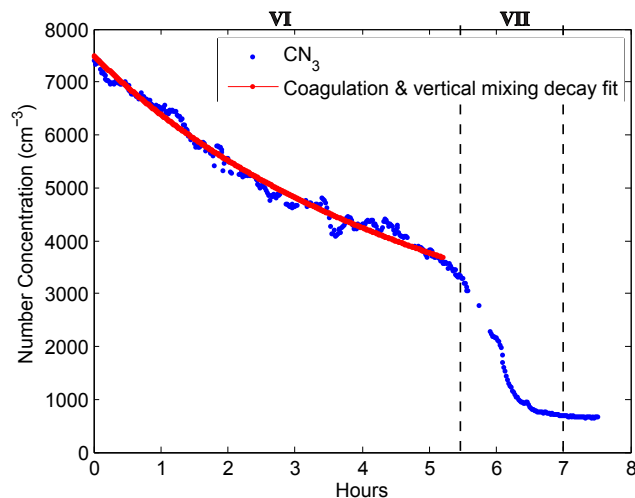


Figure 6. CN₃ data with the coagulation and vertical mixing decay fit that is found to fit for the first five hours of decay, before an air-mass change results in rapid decay to background at around 5.5 h. Zero hours corresponds to midday on the 18 October (UTC).

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

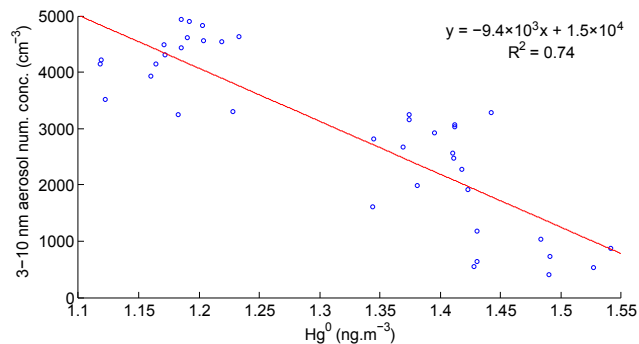


Figure 7. A negative linear correlation is found when CN_{3-10} was time lagged by two-hours relative to TGM, suggesting involvement of Hg in the nucleation chemistry. Data used show a 5 h period, starting 18 October 02:00 for TGM and 04:00 for CN data. The data gap surrounding TGM concentrations of 1.3 ng m^{-3} is a result of routine calibrations which coincided with the period of interest, and resulted in a 90 min data-gap.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

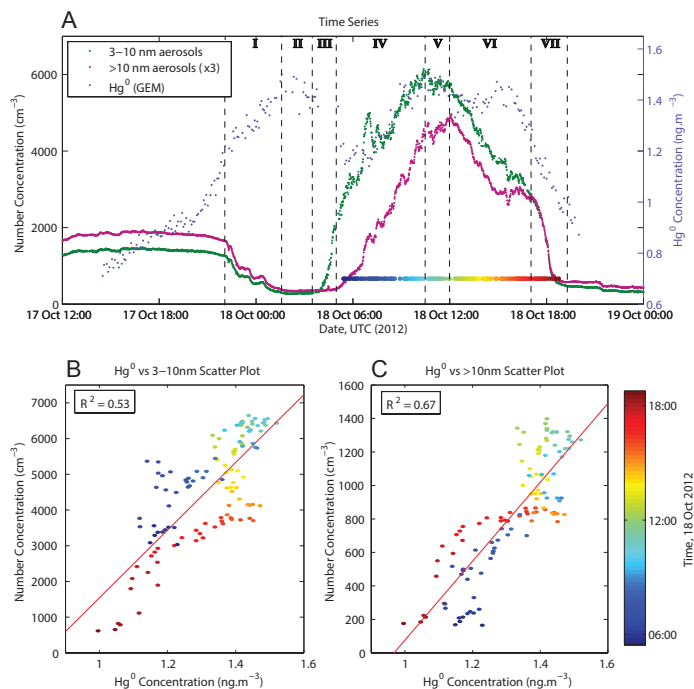


Figure 8. TGM and number concentrations of aerosols in the two measured size bins during the particle formation event. **(a)** Time series showing variations in concentrations of aerosols and TGM. **(b)** and **(c)** show correlation plots relating TGM with CN_{3-10} and CN_{10} , respectively. Both correlation plots show positive relationships, with deviations occurring primarily because of short term variations in aerosol concentrations. The correlation plots are coloured according to time of day (referenced directly in **a**) and include robust linear regression lines with corresponding R^2 values. Removal of the “wings” of plot **(b)**, a result of varying background concentrations (see text), increases the R^2 value to 0.85.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

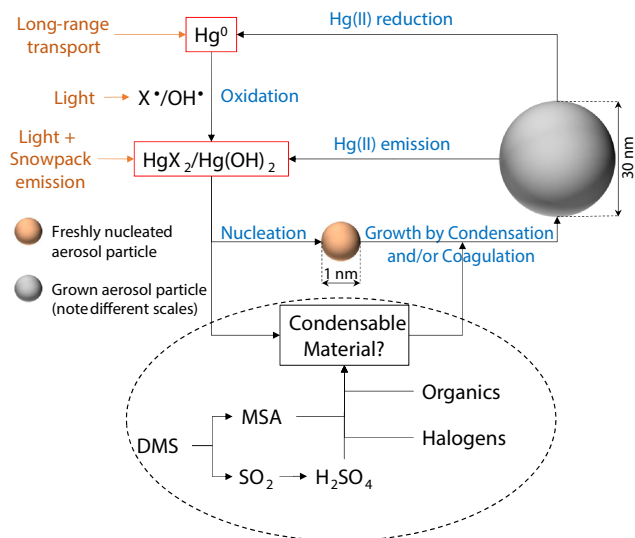


Figure 9. A suggested nucleation and growth mechanism driven by atmospheric mercury present from long-range transport. This chemistry is likely only to occur under conditions of low background aerosol surface area, otherwise these reaction products are more likely to adhere to pre-existing aerosol surfaces. Species highlighted in red constitute what is measured as TGM. Evidence suggests that no single species, DMS, organics, or halogens, can explain growth rates, however they may combine into a single reservoir of condensable material that may contribute to the growth process.

Boundary layer new particle formation over East Antarctic sea ice – possible Hg driven nucleation?

R. S. Humphries et al.

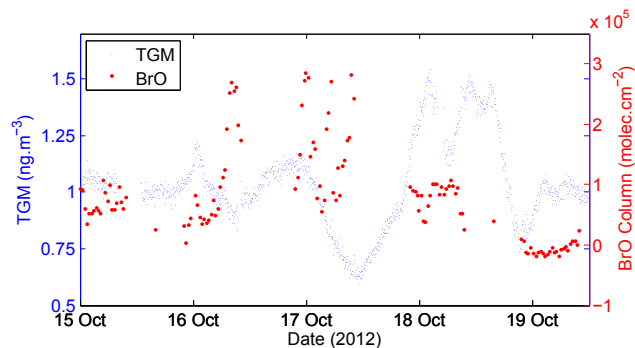


Figure A1. Total Gaseous Mercury (TGM) and column BrO from the MAX-DOAS instrument. The 17 October experiences significant concentrations of BrO, along with simultaneous depletion of TGM, just prior to the 18 October nucleation event.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)